# **OPTICAL FIBER COMMUNICATION**

#### PART III:-DIGITAL RECEIVERS

- **Digital Receivers** 
  - **Probability of error**
  - Quantum Limit
- Shot noise
- Noise Penalty
- Pre-amplifier types



Photon stream hv

Amplifier (gain M)Amplifier Bulk dark current Thermal Photon detection noise Surface leakage noise quantum noise (Poisson current • Statistical gain fluctuation) fluctuation (for avalanche photodiodes)

Photodetector

**Bias** resistor

- Analog system—Signal to rms noise ratio
- Digital system– Average error probability

#### DIGITAL RECEIVER-PHOTON DETECTION QUANTUM NOISE

- Due to random arrival rate of signal photons.
- Makes primary photocurrent a time varying Poisson Process.
- If detector illuminated by optical signal p(t), then average number of electron-hole pair generated in time τ is --

$$\overline{N} = \frac{\eta}{h\nu} \int_0^\tau P(t) dt = \frac{\eta E}{h\nu}$$

• η is detector quantum efficiency.

#### DIGITAL RECEIVER-Photon detection quantum noise

• Actual number of electron-hole pair n fluctuates from average according to Poisson distribution.

$$P_r(n) = \overline{N}^n \frac{e^{-\overline{N}}}{n!}$$

where  $P_r(n)$  is the probability that *n* electrons are emitted in an interval  $\tau$ .



• As the pulse progresses, it spreads and enters into adjacent time slots causing ISI.

 $\circ$  Major part y in desired slot while rest spreads.



- $\eta$  Quantum efficiency of photo detector
- $\circ$  C<sub>d</sub> Photodiode capacitance
- $\circ R_b$  Detector bias resistance
- ${\rm \circ}~{\rm R_a}~{\rm II}~{\rm C_a}{\rm -}~{\rm Amplifier}$  input impedance
- $\circ C_a$  Amplifier shunt capacitance
- $i_b(t)$  Thermal noise current generated by  $R_b$
- $i_a(t)$  Thermal noise current generated by  $R_a$
- $v_a(t)$  Thermal noise voltage of amplifier channel
- ${\rm \circ}$  Input voltage develops across  ${\rm R_a}$
- Two amplifier noise sources  $i_a(t)$ ,  $v_a(t)$
- One detector noise source  $i_b(t)$  due to bias resistor.

- All noises Gaussian, have flat spectral response ( white noise), uncorrelated, statically independent .
- Occurrence of one doesn't effect occurrence of other.
- Input pulse train is -

$$P(t) = \sum_{n=-\infty}^{\infty} b_n h_p(t - nT_b)$$

- $b_n$  Amplitude of n<sup>th</sup> pulse
- $h_p$  received Pulse shape
- $T_b Bit period$
- $b_n \operatorname{can} be b_{on} \operatorname{or} b_{off}$
- $h_p(t)$  normalized to have unit area.

$$\int_{-\infty}^{\infty} h_p(t) dt = 1$$

• Mean output current from detector -

$$\langle i(t) \rangle = \frac{\eta q}{h \nu} MP(t) = \mathscr{R}_0 M \sum_{n=-\infty}^{\infty} b_n h_p(t-nT_b)$$

where  $\mathscr{R}_0 = \eta q / h \nu$  is the photodiode responsivity

, NC

• Current amplified and filtered to produce mean voltage at output of equalizer -

$$\langle v_{out}(t) \rangle = A \mathscr{R}_0 MP(t) * h_B(t) * h_{eq}(t)$$
  
=  $\mathscr{R}_0 GP(t) * h_B(t) * h_{eq}(t)$   $G = AM$ 

•  $h_B(t)$  and  $h_{eq}(t)$  are impulse responses of bias and equalizer circuits.

$$h_{B}(t) = F^{-1}[H_{B}(f)] = \int_{-\infty}^{\infty} H_{B}(f)e^{j2\pi ft} df$$

$$H_{B}(f) = \frac{1}{1/R + j2\pi fC} \qquad \qquad \frac{1}{R} = \frac{1}{R_{a}} + \frac{1}{R_{b}}$$

$$C = C_{a} + C_{d}$$
The mean output voltage from equalizer is -
$$\left\langle v_{\text{out}}(t) \right\rangle = \sum_{n=-\infty}^{\infty} b_{n}h_{\text{out}}(t - nT_{b})$$

where

$$h_{\text{out}}(t) = \mathscr{R}_0 Gh_p(t) * h_B(t) * h_{\text{eq}}(t)$$

$$H_{\text{out}}(f) = \int_{-\infty}^{\infty} h_{\text{out}}(t) e^{-j2\pi ft} dt = \mathscr{R}_0 GH_p(f) H_B(f) H_{\text{eq}}(f)$$

#### PROBABILITY OF ERROR

$$BER = \frac{N_e}{N_t} = \frac{N_e}{Bt}$$

$$N_e - \text{ errors occurring in time t}$$

$$N_t - \text{ total pulses transmitted in time t}$$

$$B - \text{ bit rate} = 1/T_b$$

- Assuming the noise has Gaussian probability density function with zero mean.
- Noise voltage n(t) sampled at any arbitrary time t,
- The probability that the measured sample n(t) falls in range *n* to *n*+*dn* is -

$$f(n) dn = \frac{1}{\sqrt{2\pi\sigma^2}} e^{-n^2/2\sigma^2} dn$$

where  $\sigma^2$  is the noise variance and f(n) is the probability density function.

**PROBABILITY OF ERROR-** CASE I - 0' is being sent.

- Let transmitted pulses are '0' and '1', v<sub>th</sub> = V/2
  Transmitted = '0'
- Received r(t) = n(t) = v
- Probability that it be detected as '1' is probability that v lied between V /2and  $\infty$ .

$$P_{0}(v_{\rm th}) = \int_{V/2}^{\infty} p(y|0) \, dy = \int_{V/2}^{\infty} f_{0}(y) \, dy$$
$$= \int_{V/2}^{\infty} \frac{1}{\sqrt{2\pi\sigma^{2}}} e^{-v^{2}/2\sigma^{2}} \, dv$$

where the subscript 0 denotes the presence of a 0 bit.

**PROBABILITY OF ERROR-** CASE II - '1' IS BEING SENT.

- Transmitted = '1'
- Received r(t) = V + n(t) = v
- $\circ n(t) = v V$
- Probability that it be detected as '0' is probability that v lied between  $\infty$  and V /2.

$$f_1(v) = \frac{1}{\sqrt{2\pi\sigma^2}} e^{-(v-V)^2/2\sigma^2}$$

$$P_{1}(v_{\rm th}) = \int_{-\infty}^{V/2} p(y|1) \, dy = \int_{-\infty}^{V/2} f_{1}(v) \, dv$$



•  $P_e = a p_0(v) + b p_1(v)$ 

Assuming '0' and '1' are equiprobable, a = b = 0.5.
P<sub>o</sub> = ?

- Looking at distribution  $p_0(v)$  and  $p_1(v)$  are identical.
- Integrating double of one part -



$$P_{e} = \int_{V/2}^{\infty} \frac{1}{\sqrt{2\pi\sigma^{2}}} e^{-v^{2}/2\sigma^{2}} dv$$

•  $v/\sqrt{2} \sigma = x$ ,  $dv/\sqrt{2} \sigma = dx$ 

- Limits = ?  $P_e = ?$
- Expression not integrable. Tabulated as erfc(x) or erf(x).
- Comparing with erfc(x)
- $\operatorname{erfc}(\mathbf{x}) = 2/\sqrt{\pi}\int_{\mathbf{x}}$

$$e^{-y^2}dy$$

 $P_e = \frac{1}{2} \operatorname{erfc}(\frac{V}{2 \sigma})$ 

$$P_{e} = \frac{1}{2} \left[ 1 - \operatorname{erf} \left( \frac{V}{2\sqrt{2}\sigma} \right) \right]$$

$$\left( \frac{S}{N} \right)_{dB} = 20 \log \frac{V}{\sigma}$$
By doubling V, BER decreases by 10<sup>4</sup>.

#### QUANTUM LIMIT TO DETECTIO

$$P_r(n) = \overline{N}^n \frac{e^{-\overline{N}}}{n!}$$

- Ideal photo detector having unity quantum efficiency and no dark current.
- No e-h pair generated in absence of optical pulse.'0'
- Possible to find minimum received optical power required for specific BER performance in digital system.
- Called Quantum limit.

#### QUANTUM LIMIT TO DETECTIO

$$P_r(n) = \overline{N}^n \frac{e^{-\overline{N}}}{n!}$$

- Optical pulse of energy E falls on photo detector in time interval τ.
- During transmission signal if too low to generate any e-h pair and detected as 0.
- Then for error probability  $P_r(0)$ , there exists a minimum energy E at wavelength  $\lambda$ , to be detected as 1.
- Probability that n=0 electrons are emitted in interval τ-

$$P_r(0) = e^{-\vec{N}}$$

#### QUANTUM LIMIT TO DETECTION - PROBLEM $\overline{N} = \frac{\eta}{h\nu} \int_0^{\tau} P(t) dt = \frac{\eta E}{h\nu}$

- Digital fiber optic link operating at wavelength 850nm requires maximum BER of  $10^{-9}$ . Find quantum limit and minimum incident power P<sub>o</sub> that must fall on photo detector, to achieve this BER at data rate of 10Mbps for simple binary level signaling system. Quantum efficiency is 1.
- Solution –for maximum BER,--

$$P_r(0) = e^{-\bar{N}} = 10^{-9}$$

$$\overline{N} = 9 \ln 10 = 20.7 \approx 21.$$
$$E = 20.7 \frac{h\nu}{\eta}$$

**QUANTUM LIMIT TO DETECTION - PROBLEM** 

Minimum incident power that must fall on photo detector P<sub>o</sub> --- E= P<sub>o</sub> τ
Assuming equal number of 0 and 1, 1/ τ = B/2

$$P_0 = 20.7 \frac{hcB}{2\lambda}$$
  
=  $\frac{20.7(6.626 \times 10^{-34} \text{J} \cdot \text{s})(3.0 \times 10^8 \text{ m/s})(10 \times 10^6 \text{ bits/s})}{2(0.85 \times 10^{-6} \text{ m})}$   
= 24.2 pW

or, when the reference power level is one milliwatt,

$$P_0 = -76.2 \text{ dBm}$$

#### **RECEIVER NOISES**

• Noise voltage  $v_{\rm N}(t)$  causes  $v_{\rm out}$  to deviate from mean or average <  $v_{\rm out}$  >

$$v_{\text{out}}(t) = \langle v_{\text{out}}(t) \rangle + v_N(t)$$
$$v_N^2(t) = v_s^2(t) + v_R^2(t) + v_I^2(t) + v_E^2(t)$$

- $v_{s}(t)$  Quantum or shot noise due to random multiplied poisson nature of photocurrent  $i_{s}(t)$ .
- $v_{\rm R}(t)$  thermal noise due to bias resistor  $R_{\rm b}$ .
- $v_{\rm I}(t)$  –noise due to amplifier input noise
- $v_E(t)$  -noise due to amplifier due to  $e_a(t)$ .

• Calculating the three thermal noise currents at the output of equalizer :-

The thermal noise of the load resistor  $R_b$ 

$$\left\langle v_R^2(t) \right\rangle = \frac{4k_BT}{R_b} B_{bae} R^2 A^2$$

 $\mathbf{B}_{\mathrm{bae}}$  – Noise equivalent bandwidth of bias ckt, amplifier and equalizer

$$B_{bae} = \frac{1}{|H_B(0)H_{eq}(0)|^2} \int_0^\infty |H_B(f)H_{eq}(f)|^2 df$$
$$= \frac{1}{|H_{out}(0)/H_p(0)|^2} \int_0^\infty \left|\frac{H_{out}(f)}{H_p(f)}\right|^2 df$$

$$\left\langle v_I^2(t) \right\rangle = S_I B_{bae} R^2 A^2$$

$$\left\langle v_E^2(t) \right\rangle = S_E B_e A^2$$

where  $S_I$  is the spectral density of the amplifier input noise current source (measured in amperes squared per hertz),  $S_E$  is the spectral density of the amplifier noise voltage source (measured in volts squared per hertz), and

$$B_{e} = \frac{1}{|H_{eq}(0)|^{2}} \int_{0}^{\infty} |H_{eq}(f)|^{2} df$$
  
= 
$$\frac{R^{2}}{|H_{out}(0)/H_{p}(0)|^{2}} \int_{0}^{\infty} \left|\frac{H_{out}(f)}{H_{p}(f)} \left(\frac{1}{R} + j2\pi fC\right)\right|^{2} df$$

#### SHOT NOISES-

 $\langle v_s^2(t) \rangle = 2q \langle i_0 \rangle \langle m^2 \rangle B_{bae} R^2 A^2$ 

#### $\langle m^2 \rangle$ is the mean square avalanche gain

•Shot noise in bit period  $T_b$  is shot noise contribution from a pulse within that period as well as from all other pulses outside that period.

- •Worst case shot noise when all neighboring pulses are '1'. •Greatest ISI.
- •Hence mean unity gain photocurrent over  $T_b$  for 1 pulse -

$$\langle i_0 \rangle_1 = \sum_{n=-\infty}^{\infty} \frac{\eta q}{h\nu} b_{\text{on}} \frac{1}{T_b} \int_{-T_b/2}^{T_b/2} h_p(t - nT_b) dt$$

#### Shot Noises - '1' with all neighbor '1's

 $= \frac{\eta q}{h\nu} \frac{b_{\rm on}}{T_b} \int_{-\infty}^{\infty} h_p(t) dt$  $\frac{\eta q}{h\nu} \frac{b_{\rm on}}{T_b}$ 

SHOT NOISES - '0' WITH ALL NEIGHBOR '1'S -  

$$B_{OFF} = 0$$

$$\langle i_0 \rangle_0 = \sum_{n \neq 0} \frac{\eta q}{h\nu} b_{\text{on}} \frac{1}{T_b} \int_{-T_b/2}^{T_b/2} h_p(t - nT_b) dt$$

$$= \frac{\eta q}{h\nu} \frac{b_{\text{on}}}{T_b} \bigg[ \sum_{n = -\infty}^{\infty} \int_{-T_b/2}^{T_b/2} h_p(t - nT_b) dt - \int_{-T_b/2}^{T_b/2} h_p(t) dt \bigg]$$

$$= \frac{\eta q}{h\nu} \frac{b_{\text{on}}}{T_b} (1 - \gamma) \qquad \gamma = \int_{-T_b/2}^{T_b/2} h_p(t) dt$$

•Substitute  $\langle i_0 \rangle_0$  and  $\langle i_0 \rangle_1$  in  $v_s^2(t)$  to find worst case shot noise for pulse '1' and '0'.

#### SNR REQUIRED TO ACHIEVE MIN BER



• Assuming output voltage is approximately Gaussian.

- Mean and variance of Gaussian output for '1' and '0' are  $b_{on}$ ,  $\sigma_{on}{}^2$  and  $b_{off}$  and  $\sigma_{off}{}^2$ .
- $\bullet$  Decision threshold  $v_{th}$  set for equal error probability for '1' and '0'.

$$P_0(v_{\rm th}) = P_1(v_{\rm th}) = \frac{1}{2}P_e$$

SNR REQUIRED TO ACHIEVE MIN BER

• Error probability for '1' and '0' are

$$P_{e} = \frac{1}{\sqrt{2\pi}\sigma_{\text{off}}} \int_{v_{\text{th}}}^{\infty} \exp\left[-\frac{(v-b_{\text{off}})^{2}}{2\sigma_{\text{off}}^{2}}\right] dv$$
$$P_{e} = \frac{1}{\sqrt{2\pi}\sigma_{\text{on}}} \int_{-\infty}^{v_{\text{th}}} \exp\left[-\frac{(-v+b_{\text{on}})^{2}}{2\sigma_{\text{on}}^{2}}\right] dv$$

• Defining Q related to SNR to achieve desired min BER-

$$Q = \frac{v_{\rm th} - b_{\rm off}}{\sigma_{\rm off}} = \frac{b_{\rm on} - v_{\rm th}}{\sigma_{\rm on}}$$

• Putting  $Q/\sqrt{2} = x$ , change integral and limits.



- Relative to noise at b<sub>off</sub>, threshold v<sub>th</sub> must be ATLEAST Q standard deviation above b<sub>off</sub>.
- ${\circ}$  Or,  $v_{th}$  should be above  $b_{off}$  by rms value Q.
- $\bullet$  Relative to noise at  $b_{on},$  threshold  $v_{th}$  must not be below  $b_{on}$  by more than Q sandard deviation to have min BER

#### NOISE PENALTY IN PRACTICAL SYSTEM- POWER PENALTY

- In practical system, many factors degrade the performance.
- We assumed that
  - Optical energy of each bit is impulse response h(t).
  - Zero energy sent during '0'.
  - Receiver amplifier sharply band limited.
  - No random variation in amplitude and arrival time of bit.
- In practical systems, each violation demands increase in received signal power to ensure given error probability.
  This additional excess power ΔP required in practical

system is called power penalty – in dB

$$\Delta P = 10 \log \frac{b_{\text{on, nonimpulse}}}{b_{\text{on, impulse}}}$$

#### 1. NON-ZERO EXTINCTION RATIO

- Assumed  $b_{off} = 0$  during '0'.
- In actual system, light source biased slightly ON at all times to obtain shorter turn-on time in LED or keep it above threshold in LASER.
- Extinction ratio E is optical energy emitted in the '0' pulse to that during '1' pulse.
- $\circ \mathbf{E} = \mathbf{b}_{off} / \mathbf{b}_{on}$
- Varies between 0 and 1.
- Any dark current in photodiode appears to increase E.
- With equally probable '0' and '1', minimum received power ( sensitivity)  $P_{r min}$  is given by average energy detected per pulse times the pulse rate 1 /  $T_b$ .

1. NON-ZERO EXTINCTION RATIO



• The extinction ratio penalty i.e., the penalty in receiver sensitivity as a function of extinction ratio is-

$$y(\epsilon) = \frac{P_{r,\min}(\epsilon)}{P_{r,\min}(0)}$$

#### 2. FINITE PULSE WIDTH AND TIMING JITTER

- Received optical pulse has a finite pulse width.• Some timing jitter is present.
- Hence noise penalty is required for—
  - Non-optical filtering is needed to provide equalization against pulse distortion or to minimize ISI.
  - Some ISI remains and degrades SNR.
- To calculate magnitude of these effects, it is necessary to define
  - Shape of received pulse.
  - The distribution of the jitter

• We deal with only the former.

#### 2. FINITE PULSE WIDTH AND TIMING JITTER



Power penalty Vs ρ/T for Gaussian shaped pulse shown.
ρ is rms width of the pulse (due to changes in pulse shape.) T is basic pulse width.

#### 2. FINITE PULSE WIDTH AND TIMING JITTER

- It demonstrates possible trade-off between bit rate and signal power.
- Relates effects of fiber attenuation and fiber disperssion.
- Power penalty < 1db if  $\rho$  remains less than T/5.
- Result independent of pulse width.
- But if  $\rho > 1$ dB, PP increases sharply.
- It becomes more sensitive to pulse shape.
- In practice, system is
  - either limited by fiber dispersion (BW Limited)
  - or by fiber attenuation(Power Limited) .
  - Possible trade-off between two is quite small.

#### PREAMPLIFIER TYPES

- Sensitivity and bandwidth of a receiver are effected by noise sources at the front end, i.e. at preamplifier.
- Preamplifier should give maximum receiver sensitivity with desired bandwidth.
- Three main types. But intermediate types can also be used.

#### PREAMPLIFIER TYPES – LOW IMPEDANCE (LZ)

- Simplest, but not optimum design.
- Photodiode operates into a low impedance amplifier (appox. $50\Omega$ )
- Bias or load resister R<sub>b</sub> used to match amplifier impedance by suppressing standing waves and give uniform frequency response.
- R<sub>b</sub> with amplifier capacitance gives BW equal to or greater than signal BW.
- LZ amplifier can operate over a wide BW.
- ${\rm \circ}$  Gives low resistivity as a small voltage develops across amplifier input and  $R_{\rm b}$  .
- Hence used only for short distance applications where sensitivity not of concern.



## PREAMPLIFIER TYPES – HIGH IMPEDANCE (HZ)

- Noise reduced by reducing input capacitance.
- By selecting.
  - Low capacitance, high frequency devices.
  - Detector of low dark current.
  - Bias resistor of minimum thermal noise.
- Thermal noise can be reduced by high impedance amplifier with large photo detector R<sub>b</sub>, hence HZ amp<sup>r</sup>.
- But causes large RC time constant low front-end BW.
- Signal gets integrated.
- Equalization technique required. (Differentiator).
- Integrator-differentiator is core of HZ design.
- Gives low noise but low dynamic range of signal.





HIGH IMPEDANCE- NOISE OUTPUT  $V_{\rm N}^2 = \int_{\Delta f} |G(f)|^2 (V_{\rm A}^*)^2 \, \mathrm{d}f + \int_{\Delta f} \frac{|G(f)|^2 R^2 (I_{\rm T}^*)^2 \, \mathrm{d}f}{|1 + j2\pi f C R|^2}$  $V_{\rm N}^2 = G_0^2 \int_0^{\Delta y} \left\{ (1 + 4\pi^2 f^2 C^2 R^2) (V_{\rm A}^*)^2 + R^2 (I_{\rm T}^*)^2 \right\} df$  $V_{\rm N} = \left[ \left\{ 1^{*} + (4\pi^{2}/3)(\Delta f)^{2}C^{2}R^{2} \right\} (V_{\rm A}^{*})^{2} + R^{2}(I_{\rm T}^{*})^{2} \right]^{1/2} (\Delta f)^{1/2}$  $K = \frac{V_{\text{out}}}{|V_{\text{N}}|} = \frac{MIR}{\{(1 + \frac{4}{3}\pi^{2}(\Delta f)^{2}C^{2}R^{2})(V_{\text{A}}^{*})^{2} + R^{2}(I_{\text{T}}^{*})^{2}\}^{1/2}(\Delta f)^{1/2}}$  $K = \frac{1}{\left\{\frac{(V_A^*)^2}{M^2} \left(\frac{1}{R^2} + \frac{4\pi^2}{3}(\Delta f)^2 C^2\right) + 2eIF + \frac{4kT}{M^2R} + \frac{(I_A^*)^2}{M^2}\right\}^{1/2} (\Delta f)^{1/2}}$ (b) (c) (d) (e)

#### HIGH IMPEDANCE (HZ)- ANALYSIS

- SNR K can be improved by increasing M until shot noise term (iii) increases by F(M).
- F(M) increases with M, gets comparable to other terms. Has optimum value for M.
- K improves by increasing front end R till (i) and (iv) are significant. But increases RC. Hence requires more equalization and low C.
- If equalization required, (ii) dominates at high freq. Noise increases as C<sup>2</sup>. Hence low C required.
- Shot noise (iii) depends on input signal level.
- Assumption of all noises statistical not true in reality.
- All noises including F(M) not purely Gaussian.

#### HIGH IMPEDANCE (HZ)- ANALYSIS

- Hence Actual SNR may be lesser.
- Two limitations due to Integrator-differentiator
- Equalization required for broadband applications.
- limited dynamic range.

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- High gain and high impedance amplifier with feedback.
- Gives low noise and large dynamic range.
- $\circ 1/R = 1/R_a + 1/R_b + 1/R_f$
- $\circ C = C_a + C_d$

TRANS IMPEDANCE – SIGNAL POWER 
$$V_{in} = -V/A$$
  
 $MI + \frac{V - V_{in}}{T_F} = V_{in} \left(\frac{1}{R} + j2\pi fC\right)$   
 $MI = -V\left(\frac{1}{R_F} + \frac{1}{AR_F} + \frac{1}{AR} + \frac{j2\pi fC}{A}\right)$   
 $V = \frac{-R_F MI}{1 + \frac{1}{A} + \frac{R_F}{AR} + j\frac{2\pi fCR_F}{A}} = \frac{-R_F MI/(1 + 1/A + R_F/AR)}{[1 + j2\pi fCR_F/(1 + R_F/R + A)]}$   
 $A \gg 1 + R_F/R$   
 $V \simeq \frac{-R_F MI}{(1 + j2\pi fCR_F/A)}$ 



TRANS IMPEDANCE – NOISE POWER  

$$(V_F^*)^2 = 4kTR_F$$

$$\left(V_A^* + \frac{V_{N_1}^*}{(-A)}\right)\left(\frac{1}{R} + j2\pi fC\right) = \left(V_{N_1}^* - V_A^* - \frac{V_{N_1}^*}{(-A)}\right)\frac{1}{R_F}$$

$$V_{N_1}^* = V_A^*R_F\left(\frac{1}{R_F} + \frac{1}{R} + j2\pi fC\right) / \left(1 + \frac{1}{A} + \frac{R_F}{AR} + \frac{j2\pi fCR_F}{A}\right)$$

$$\simeq V_A^*\left(1 + \frac{R_F}{R} + j2\pi fCR_F\right)$$

.

#### TRANS IMPEDANCE – NOISE POWER

For the current source

$$\left(V_{N2}^{*} - \frac{V_{N2}^{*}}{(-A)}\right) \frac{1}{R_{F}} + I_{T}^{*} = \frac{V_{N2}^{*}}{(-A)} \left(\frac{1}{R} + j2\pi fC\right)$$
$$V_{N2}^{*} \left(\frac{1}{R_{F}} + \frac{1}{AR_{F}} + \frac{1}{AR_{F}} + \frac{1}{AR} + \frac{j2\pi fC}{A}\right) = -I_{T}^{*}$$

 $\therefore V_{\rm N2}^* \simeq -R_{\rm F}/I_{\rm T}^*$ 

#### TRANS IMPEDANCE – NOISE POWER

#### For the feedback source

$$V_{N3}^* = V_F^* + \frac{V_{N3}^*}{(-A)} = \frac{V_F^*}{(1+1/A)} \simeq V_F^*$$

 $(V_{\rm N}^{*})^{2} = (V_{\rm N1}^{*})^{2} + (V_{\rm N2}^{*})^{2} + (V_{\rm N3}^{*})^{2}$  $= (V_{\rm A}^{*})^{2} \left\{ \left(1 + \frac{R_{\rm F}}{R}\right)^{2} + 4\pi^{2}f^{2}C^{2}R_{\rm F}^{2} \right\} + R_{\rm F}^{2}(I_{\rm T}^{*})^{2} + (V_{\rm F}^{*})^{2}$ 

#### TRANS IMPEDANCE – NOISE POWER

### PREAMPLIFIER TYPES– TRANS IMPEDANCE - ANALYSIS

- R and  $R_f$  can be increased to reduce SNR, (i) and (iv) without equalization provided A>> $2\pi C R_f B$ .
- It has wide dynamic range.
- Output resistance is small so that amplifier is less susceptible to pickup noise.
- Transfer characteristic is actually trans impedance feedback resister. Hence amplifier stable and easily controlled.
- Although TZ amp<sup>r</sup> is less sensitive than HZ amp<sup>r</sup> as  $S/N_{TZ} > S/N_{HZ}$ , the difference is usually only 2 to 3 dB for most practical wideband design.